

Phys 262 - PARTICLE IN A BOX, CHAPTER 40

IN QUANTUM MECHANICS THE PROBABILITY DISTRIBUTION FUNCTION IS GIVEN BY $|\psi|^2 = \psi\psi^*$

FOR A STATIONARY STATE, $\psi = \phi e^{-iEt/\hbar}$

$$\psi\psi^* = \phi e^{-iEt/\hbar} \phi^* e^{iEt/\hbar} = |\phi|^2$$

AVERAGES ARE FOUND BY INTEGRATING THE DISTRIBUTION FUNCTION:

$$\bar{x} = \int x |\psi|^2 dx = \int x |\phi|^2 dx$$

IN QUANTUM MECHANICS, WE CAN ALSO FIND THE AVERAGE MOMENTUM, \bar{p} .

CONSIDER THE FREE PARTICLE, $U = 0 \Rightarrow \phi = Ae^{ikx} + Be^{-ikx}$

WHERE $p = \hbar k$.

$$\begin{aligned} \text{NOTICE THAT } \frac{\partial \phi}{\partial x} &= ik Ae^{ikx} - ik Be^{-ikx} = \frac{i p}{\hbar} Ae^{ikx} + \frac{i p}{\hbar} Be^{-ikx} \\ &= \frac{i p}{\hbar} \phi \end{aligned}$$

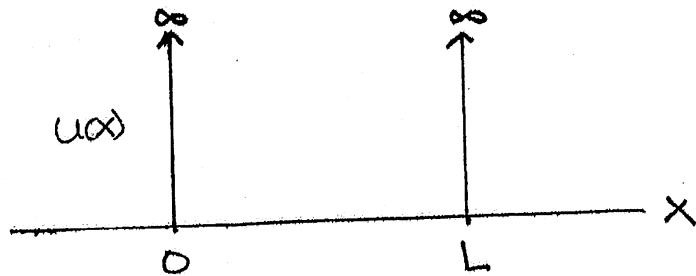
THE NEGATIVE SIGN GOES AWAY BECAUSE e^{-ikx} IS A PLANE WAVE PROPAGATING TO THE LEFT, I.E., WITH MOMENTUM $p = -\hbar k$. WE'LL ASSUME THIS IS TRUE IN GENERAL

$$\frac{\partial \phi}{\partial x} = \frac{i p}{\hbar} \phi \Rightarrow p \phi = \hbar \frac{\partial \phi}{\partial x} = \frac{i}{i} \hbar \frac{\partial \phi}{\partial x} = \frac{i \hbar}{i^2} \frac{\partial \phi}{\partial x} = -i \hbar \frac{\partial \phi}{\partial x}$$

$$\bar{p} = \int p |\psi|^2 dx = \int p |\phi|^2 dx = \int p \phi \phi^* dx = \int \phi^* p \phi dx$$

$$\Rightarrow \boxed{\bar{p} = \int \Psi^* (-i\hbar) \frac{\partial \Psi}{\partial x} dx}$$

PARTICLE IN A BOX — Box \Rightarrow REGION OF LENGTH L WHOSE SIDES ARE INFINITELY STRONG.



$$U(0) = U(L) = \infty$$

\Rightarrow PROBABILITY TO BE OUTSIDE THE BOX IS ZERO $\Rightarrow \Psi(0) = \Psi(L) = 0$.

INSIDE THE BOX, $U(x) = 0 \Rightarrow \Psi = Ae^{ikx} + Be^{-ikx}$

NEED TO MATCH THE BOUNDARY CONDITIONS TO FIND A AND B

BC'S ARE $\Psi(0) = 0, \Psi(L) = 0$

$$\Psi(0) = Ae^{ik \cdot 0} + Be^{-ik \cdot 0} = A + B \Rightarrow A + B = 0 \Rightarrow A = -B$$

$$\Rightarrow \Psi = Ae^{ikx} - Ae^{-ikx} = A(e^{ikx} - e^{-ikx}) = A(2i \sin(kx)) = 2iA \sin(kx)$$

$$\Rightarrow \Psi = C \sin(kx)$$

$$\Psi(L) = C \sin(kL) = 0 \Rightarrow kL = n\pi \quad n=1, 2, 3, \dots \quad (\text{CAN'T DO } n=0 \text{ BECAUSE } L \neq 0 \text{ NOR DOES } k=0)$$

$$p = \hbar k, \quad U = 0, \quad E = K + U \Rightarrow E = K = \frac{p^2}{2m}$$

$$\Rightarrow E = \frac{(\hbar k)^2}{2m} = \frac{\hbar^2}{2m} \left(\frac{n\pi}{L} \right)^2 \Rightarrow \boxed{E = \frac{\pi^2 \hbar^2}{2mL^2} n^2}$$

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WE FIND THE C CONSTANT FROM NORMALIZATION

$\int_{-\infty}^{\infty} |\Phi|^2 dx = 1$. THE PROBABILITY TO BE OUTSIDE THE BOX ($x < 0, x > L$)

IS 0 $\Rightarrow \Phi = 0$ FOR $x < 0$ AND $x > L \Rightarrow$

$$\int_{-\infty}^{\infty} |\Phi|^2 dx = \int_0^L |\Phi|^2 dx = \int_0^L C^2 \sin^2 kx dx = C^2 \int_0^L \sin^2 kx dx$$

$$k = \frac{n\pi}{L} \Rightarrow C^2 \int_0^L \sin^2\left(\frac{n\pi x}{L}\right) dx = C^2 \int_0^L \frac{1}{2} (1 - \cos\left(\frac{2n\pi x}{L}\right)) dx$$

$$= \frac{C^2}{2} \left(L - \frac{L}{2n\pi} \sin\left(\frac{2n\pi x}{L}\right) \Big|_0^L \right) = \frac{C^2}{2} (L - 0) = \frac{C^2 L}{2} \Rightarrow \frac{C^2 L}{2} = 1$$

$$\Rightarrow C^2 = \frac{2}{L} \Rightarrow C = \sqrt{\frac{2}{L}}$$

$$\Phi_n = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right) \text{ WITH } E_n = \frac{\pi^2 \hbar^2}{2mL^2} n^2$$

STATIONARY STATES/ENERGIES FOR PARTICLE IN A BOX.

EXAMPLE AN ELECTRON IS CONFINED TO A BOX OF LENGTH 10^{-10} m. WHAT ARE THE POSSIBLE ELECTRON ENERGIES?

$$E_n = \frac{\pi^2 \hbar^2}{2mL^2} n^2 = \frac{\pi^2 (1.06 \times 10^{-34} \text{ J}\cdot\text{s})^2}{2(9.11 \times 10^{-31} \text{ kg})(1 \times 10^{-10} \text{ m})^2} n^2 = 6.09 \times 10^{-18} \text{ J } n^2 = 38 \text{ eV } n^2$$

$$\text{UNIT: } \frac{\text{J}^2 \text{s}^2}{\text{kg m}^2} = \text{J} \left(\frac{\text{J s}^2}{\text{kg m}^2} \right) = \text{J} \left(\frac{\text{kg m}^2 / \text{s}^2 \cdot \text{s}^2}{\text{kg m}^2} \right) = \text{J} \rightarrow \text{NOTE: YOU HAVE TO USE } \hbar = 1.06 \times 10^{-34} \text{ J}\cdot\text{s} \text{ WHEN USING kg AND METERS.}$$

for $n=1$, WHAT IS THE AVERAGE VALUE OF POSITION AND MOMENTUM?

$$\Phi_1 = \sqrt{\frac{2}{L}} \sin\left(\frac{\pi x}{L}\right) \quad (\text{for } 0 < x < L)$$

$$\bar{X} = \int_0^L x |\Phi_1|^2 dx = \int_0^L x \left(\frac{2}{L}\right) \sin^2\left(\frac{\pi x}{L}\right) dx = \frac{2}{L} \int_0^L x \sin^2\left(\frac{\pi x}{L}\right) dx$$

$$= \frac{2}{L} \int_0^L x \cdot \frac{1}{2} (1 - \cos \frac{2\pi x}{L}) dx = \frac{2}{L} \cdot \frac{1}{2} \int_0^L x (1 - \cos \frac{2\pi x}{L}) dx$$

$$= \frac{1}{L} \int_0^L x dx - \underbrace{x \cos \frac{2\pi x}{L}}_{\text{INTEGRATION BY PARTS}} dx = \frac{1}{L} \left(\frac{1}{2} x^2 \Big|_0^L - \phi \right) = \frac{1}{L} \cdot \frac{1}{2} L^2 = \frac{1}{2} L$$

INTEGRATION
BY PARTS

I'll LET YOU VERIFY

$$\bar{P} = \int_0^L \Phi^* (-i\hbar) \frac{\partial \Phi}{\partial x} = \int_0^L \Phi^* (-i\hbar) \frac{\partial}{\partial x} \left(\sqrt{\frac{2}{L}} \sin\left(\frac{\pi x}{L}\right) \right) dx$$

$$= \int_0^L \Phi^* (-i\hbar) \sqrt{\frac{2}{L}} \frac{\pi}{L} \cos\left(\frac{\pi x}{L}\right) dx = \int_0^L \sqrt{\frac{2}{L}} \sin\left(\frac{\pi x}{L}\right) (-i\hbar) \sqrt{\frac{2}{L}} \frac{\pi}{L} \cos\left(\frac{\pi x}{L}\right) dx$$

$$= \frac{2}{L} (-i\hbar) \frac{\pi}{L} \int_0^L \sin\left(\frac{\pi x}{L}\right) \cos\left(\frac{\pi x}{L}\right) dx = \frac{2}{L} (-i\hbar) \frac{1}{2} \sin^2\left(\frac{\pi x}{L}\right) \Big|_0^L = \phi$$

SO ON AVERAGE THE PARTICLE IS AT THE MIDDLE OF THE BOX
AND HAS ZERO MOMENTUM.