

## Free Particle

So let's begin our exploration of solutions to the Time-Independent Schrödinger's equation with the simplest situation. Let the time independent potential be a constant with respect to position. Then, we know that the value of the constant has no impact on the behavior of the particle, so we set the value to zero,  $V(x) = V = 0$ . Then, the TISE becomes:

$$-\frac{\hbar^2}{2m} \frac{d^2\psi(x)}{dx^2} = E\psi(x)$$

The general solution to this differential equation (which can be found by integrating twice) is:

$$\psi(x) = Ae^{ikx} + Be^{-ikx}, \quad k \equiv \frac{\sqrt{2mE}}{\hbar}$$

and we recover the free particle wave function that we postulated earlier:

$$\Psi(x, t) = \psi(x)\varphi(t) = (Ae^{ikx} + Be^{-ikx})e^{-i\omega t} = Ae^{i(kx - \omega t)} + Be^{-i(kx - \omega t)}, \quad \omega = \frac{E}{\hbar}$$

Here,  $k$  can take on any positive value (depending upon the value of  $E$ ). We can let  $k$  take on both positive and negative values, where  $k > 0$  corresponds to waves traveling to the right, and  $k < 0$ , to the left. We then have:

$$\Psi(x, t) = Ae^{i(kx - \omega t)}$$

as before, or, in terms of  $k$  alone:

$$\Psi_k(x, t) = Ae^{i\left(kx - \frac{\hbar k^2}{2m}t\right)}$$

Now, this wave function should really bother you. It doesn't? OK, then, normalize it:

$$\begin{aligned} \int_{-\infty}^{\infty} \Psi_k^*(x, t)\Psi_k(x, t)dx &= \int_{-\infty}^{\infty} A^2 e^{i\left(kx - \frac{\hbar k^2}{2m}t\right)} e^{-i\left(kx - \frac{\hbar k^2}{2m}t\right)} dx \\ &= A^2 \int_{-\infty}^{\infty} dx = A^2(\infty) \end{aligned}$$

So that these wave functions cannot represent real particles, at least with single values of the wave number  $k$ , since the probability to find the particle extends to plus/minus infinity. We can address this by remembering that the general solution is a linear combination of the separable solutions. Then, as before, we can add wave functions with

a distribution of wave numbers together to form a wave packet, such that the total wave function vanishes at infinity and is therefore normalizable. Then, as before, we have for the general solution:

$$\Psi(x, t) = \int_{-\infty}^{\infty} \phi(k) \Psi_k(x, t) dk,$$

where  $\phi(k)$  now represents the constants  $c_n$  that we saw in the general solution, but is a continuous set, rather than having discrete indices so we integrate rather than sum. We find the set of constants in the usual way:

$$\Psi(x, 0) = \int_{-\infty}^{\infty} \phi(k) \psi_k(x) dk,$$

and the  $\phi(k)$  can be found by the inverse Fourier transform.