

Formula Sheet

PHYC/ECE464

Hermite-Gaussian Beams

$$\frac{E(x,y,z)}{E_0} = H_m\left(\frac{\sqrt{2}x}{w(z)}\right) H_p\left(\frac{\sqrt{2}y}{w(z)}\right) e^{-\frac{x^2+y^2}{w(z)^2}} \frac{w_0}{w(z)} e^{-i\frac{kr^2}{2q(z)}} \times \exp\left(-i\left[kz - (1+m+p)\tan^{-1}\left(\frac{z}{z_0}\right)\right]\right)$$

Low-order HG modes $\begin{cases} H_0(u) = 1; H_1(u) = 2u \\ H_2(u) = 2(2u^2 - 1) \end{cases}$

$$\frac{1}{q(z)} = \frac{1}{R(z)} - i\frac{\lambda_0}{\pi n w^2(z)}, \quad w^2(z) = w_0^2\left(1 + \frac{z^2}{z_0^2}\right), \quad R(z) = z\left(1 + \frac{z_0^2}{z^2}\right), \quad z_0 = \frac{\pi n w_0^2}{\lambda_0}, \quad k = n\frac{\omega}{c} = \frac{2\pi n}{\lambda_0}, \quad \theta_0 = \frac{2\lambda}{\pi w_0}$$

Optics

Irradiance $I = \langle S \rangle = \frac{n c \epsilon_0}{2} E_0^2$

Fresnel coefficients:

$$r_{\parallel} = \frac{n_t \cos(\theta_i) - n_i \cos(\theta_t)}{n_t \cos(\theta_i) + n_i \cos(\theta_t)} = \frac{\tan(\theta_i - \theta_t)}{\tan(\theta_i + \theta_t)}$$

$$r_{\perp} = \frac{n_i \cos(\theta_i) - n_t \cos(\theta_t)}{n_i \cos(\theta_i) + n_t \cos(\theta_t)} = \frac{\sin(\theta_i - \theta_t)}{\sin(\theta_i + \theta_t)}$$

Intensity reflectivity: $R = |r|^2$

Snell Law: $n_1 \sin(\theta_1) = n_2 \sin(\theta_2)$

Brewster and critical angles 1->2:

$$\tan(\theta_B) = (n_2/n_1)$$

$$\sin(\theta_c) = (n_1/n_2)$$

n complex: $n \rightarrow \tilde{n} = n + ik$

Lens transformation Gaussian beam:

$$\frac{1}{R_2} = \frac{1}{R_1} - \frac{1}{f}$$

Lens-maker s' formula:

$$\frac{1}{f} = (n - 1) \left(\frac{1}{R_1} - \frac{1}{R_2} \right)$$

Fabry-Perot Transmission and Reflection (with gain/loss)

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| $T(\theta, G_0) = \frac{G_0(1-R_1)(1-R_2)}{(1-G_0\sqrt{R_1R_2})^2 + 4G_0\sqrt{R_1R_2}\sin^2(\theta)}$ $R(\theta, G_0) = \frac{(\sqrt{R_1} - \sqrt{R_2})^2 + 4G_0\sqrt{R_1R_2}\sin^2(\theta)}{(1-G_0\sqrt{R_1R_2})^2 + 4G_0\sqrt{R_1R_2}\sin^2(\theta)}$ $2\Delta\theta_{1/2} = \frac{1 - G_0\sqrt{R_1R_2}}{G_0^{1/2}\sqrt{R_1R_2}}$ $\theta = kd = \frac{\omega nd}{c}$ | <p>Finesse:</p> $F = \frac{\pi^4 \sqrt{R_1 R_2}}{1 - \sqrt{R_1 R_2}} = \frac{\Delta\nu_{FSR}}{\Delta\nu_{1/2}}$ <p>FSR: $\Delta\nu_{FSR} = \frac{c}{2nd} = \frac{1}{\tau_{RT}}$</p> <p>Photon: $\tau = \frac{\tau_{RT}}{1 - R_1 R_2} \approx \frac{1}{2\pi\Delta\nu_{1/2}}$</p> <p>Resonance condition: $RTPC = q2\pi$</p> |
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ABCD matrices $\begin{pmatrix} A & B \\ C & D \end{pmatrix}$, $\det(.)=1$; $\begin{pmatrix} r_2 \\ r_2' \end{pmatrix} = \begin{pmatrix} A & B \\ C & D \end{pmatrix} \begin{pmatrix} r_1 \\ r_1' \end{pmatrix}$, Cavity s-round trip $r_s = r_{max} \sin(s\theta + \alpha)$, $-1 < \cos\theta = (A+D)/2 < 1$, $\frac{1}{q} = -\frac{A-D}{2B} - i\sqrt{1 - \left(\frac{A+D}{2}\right)^2}$ **q parameter at a cavity**

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| Free Space of length d $\begin{pmatrix} 1 & d \\ 0 & 1 \end{pmatrix}$ | Dielectric interface 1->2 $\begin{pmatrix} 1 & 0 \\ 0 & n_1/n_2 \end{pmatrix}$ | ABCD law for Gaussian beams $q_2 = \frac{Aq_1+B}{Cq_1+D}$, $q(z) = z + iz_0$ Stability: $-1 < \cos\theta = (A+D)/2 < 1$; |
| Medium of length d , index n immersed in vacuum $n_1=1$ $\begin{pmatrix} 1 & d/n \\ 0 & 1 \end{pmatrix}$ | Thin lens $\begin{pmatrix} 1 & 0 \\ -1/f & 1 \end{pmatrix}$ | Gaussian Pulse propagation $ng = n + v \frac{dn}{dv}$; $Vg = c/ng$ $\tau_p^2(z) = \tau_{p0}^2 \left(1 + \frac{z^2}{L_{GVD}^2}\right)$, $L_{GVD} = \frac{\tau_{p0}^2}{ \beta_2 }$ GVD: $\beta_2 = \frac{\lambda^3}{2\pi c^2} \frac{d^2 n}{d\lambda^2} = \frac{\lambda^2}{2\pi c} D$ |
| Curved mirror (R) $(R>0) \rightarrow$ $\begin{pmatrix} 1 & 0 \\ -2/R & 1 \end{pmatrix}$ $(R<0) \rightarrow$ | Spherical dielectric interface $\begin{pmatrix} 1 & 0 \\ ((1 - n_1/n_2)/R & n_1/n_2) \end{pmatrix}$ | Photon density $\frac{N_p}{V} = \frac{I}{h\nu c/n_g}$ |

Gain in a two-level system: $\gamma(\nu) = \sigma(\nu) \left[N_2 - \frac{g_2}{g_1} N_1 \right]$. Gain cross section: $\sigma(\nu) = A_{21} \frac{\lambda^2}{8\pi n^2} g(\nu)$

Lineshape normalization $\int g(\nu) d\nu = 1$ Beer's law: $\frac{1}{I} \frac{dI}{dz} = -\alpha(I) + \gamma(I)$

Gain/ absorption saturation, homogeneously- broadened system:

$$\gamma(I) = \frac{\gamma_0}{1+I/I_S} \quad \alpha(I) = \frac{\alpha_0}{1+I/I_S} \quad I_S = \frac{h\nu}{\sigma(\nu)\tau_2}$$

Einstein's relations and black body: $\frac{A_{21}}{B_{21}} = \frac{8\pi n^3 h\nu^3}{c^3}$, $g_2 B_{21} = g_1 B_{12}$; $\rho(\nu) = \frac{8\pi n^3 \nu^2}{c^3}$; $\frac{N_2}{N_1} = \frac{g_2}{g_1} e^{-(E_2-E_1)/kT}$

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| <p><i>Lorentzian lineshape</i></p> $g(\nu) = \frac{\Delta\nu_h/2\pi}{(\nu-\nu_0)^2+(\Delta\nu_h/2)^2}$ <p><i>Normalized Lorentzian</i></p> $\tilde{g}(\nu) = \frac{g(\nu)}{g(\nu_0)} = \frac{(\Delta\nu_h/2)^2}{(\nu-\nu_0)^2+(\Delta\nu_h/2)^2}$ | <p><i>Doppler broadened lineshape</i></p> $g(\nu) = \left(\frac{4\ln 2}{\pi}\right)^{1/2} \frac{1}{\Delta\nu_D} \exp\left[-4\ln 2 \left(\frac{\nu-\nu_0}{\Delta\nu_D}\right)^2\right] \quad \Delta\nu_D = \left(\frac{8kT\ln 2}{Mc^2}\right)^{1/2} \nu_0$ <p><i>Saturated amplifier</i></p> $\ln(G) + \frac{I_1}{I_S} (G - 1) = \gamma_0(\nu) l_g$ |
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$$\frac{dN_p}{dt} = \frac{G^2 S - 1}{\tau_{RT}} N_p + N_2 c \sigma \quad (\text{Photon number dynamics due to stimulated and spontaneous emission})$$

S (survival factor) = $R_1 R_2$ for a simple two mirror linear cavity; $G^2 = \text{r.t. gain} (\exp(2 \gamma L_g))$
 Threshold condition: $SG^2 = 1$ (linear cavity), $SG = 1$ (ring cavity)

Schawlow-Townes limit for laser linewidth: $\Delta\nu_{OSC} \approx 2\pi \frac{h\nu}{P_{OUT}} (\Delta\nu_{1/2})^2$

At steady-state: $\gamma = \gamma_{th} = \frac{\gamma_0}{1+I/I_S}$ (for homogeneously broadened)

Inside the gain medium: $I \approx I^+ + I^- \approx 2I^+$ for a high-Q linear (standing-wave) or bidirectional ring cavity, $I \approx I^+$ for a unidirectional ring cavity.

$I_{out} = T_a \cdot T_2 I^+$ (T_2 is the output coupling transmission and $T_a \dots$ are the transmission of other optical surfaces in the path)

Self consistency in ring laser provides intensity I_2 at the end of the gain medium $I_2 = I_s \frac{\gamma_0 l_g - \ln(1/S)}{1-S}$

Q-Switching and Gain-Switching: $\Delta t_p \approx \tau_p$ (cavity photon lifetime)

Power output $P_{out} = h\nu \frac{N_p}{\tau_p} \eta_{cpl}$ with η_{cpl} the coupling efficiency

With Photon number $N_p(f) = \frac{1}{2} (n_{th} \ln(n_f/n_i) + (n_i - n_f))$ with n_i the initial inversion of atoms.

Modelocking: Repetition Rate = $1/T_{rt} = 2L_n g/c$ (linear cavity), Pulsewidth: $\Delta t_p \approx 1/\Delta\nu$

Threshold current density in a diode laser: $J_{th} = eN_{eh}^{th} d/\tau_r$

Physical Constants

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| $c = 2.998 \times 10^8 \text{ m} \cdot \text{s}^{-1}$ | $h = 6.63 \times 10^{-34} \text{ J} \cdot \text{s}$ | $e = 1.6 \times 10^{-19} \text{ C}$ |
| $k_B = 1.380 \times 10^{-23} \text{ J} \cdot \text{K}^{-1}$ | $\epsilon_0 = 8.85 \times 10^{-12} \text{ F/m}$ | $m_e = 9.1 \times 10^{-31} \text{ kg}$ |

1G = 10^{-4} T, 1 eV = 1.602×10^{-19} J, 1 dyne = 10^{-5} N, 1 erg = 10^{-7} J